

# SOLAR SPECTRUM FORMATION: EXAMPLES

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The clickers below link to specific pages in this pdf file. It is a snapshot selection from my example displays made on October 4 2016 by retaining only those that concern the ALC7 model atmosphere and non-equilibrium  $H\alpha$  formation. Many are hyperlinked in my publication “*Solar  $H\alpha$  features with hot onsets III. Long fibrils in  $Ly\alpha$  and with ALMA*” accepted for Astronomy & Astrophysics on that date (DOI <http://dx.doi.org/10.1051/0004-6361/201629238>).

My current collection of lecture displays resides at:

<http://www.staff.science.uu.nl/~rutte101/Lectures.html>

**ALC7 model:**    model    hydrogen    versus FCHHT-B    strong lines

**explanation of line plots:**    pops plot    BSJ plot    profile plot

**lines from ALC7:**    Mg I 4571    Fe I 6302    Mg I b<sub>2</sub>    Na I D<sub>1</sub>    Ba II 4554  
Ca II 8542 Å    Ca II K    Mg II k    Ly $\alpha$     H $\alpha$     H $\alpha$  S

**dynamic chromosphere:**    non-E H $\alpha$     aureole boosting    H $\alpha$  extinction

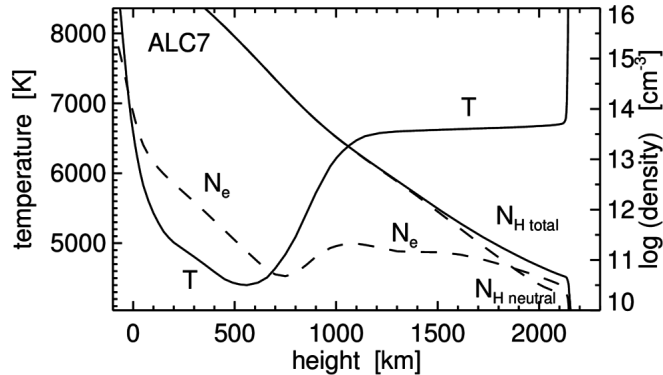
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## 2 ALC7 ATMOSPHERE

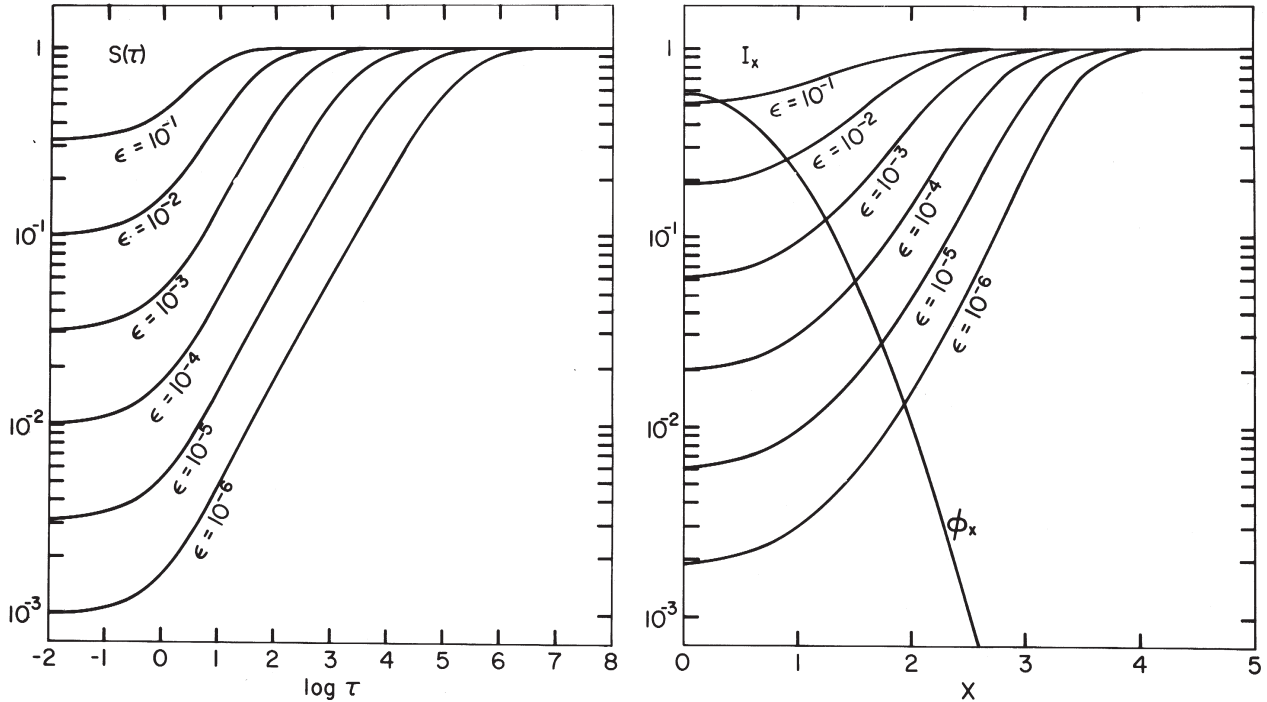
Avrett & Loeser 2008ApJS..175..229A



- unrealistic plane-parallel static computational star with solar-like average spectrum
  - exemplary in obeying all equations in my RT courses: understandable line formation
- best-fit temperature: near-RE in photosphere, shock-dominated in chromosphere
  - slope in upper photosphere depends on NLTE ultraviolet line haze
- total hydrogen density: exponential decay
  - turbulent pressure added to gain scale height and chromospheric extent
- low electron density in photosphere and temperature minimum
  - from ionization of donor-elements Si, Fe, Mg, Al with  $10^{-4}$  relative abundance
- increasing hydrogen ionization across chromosphere
  - electron density reaches proton density at its top
- near-isothermal near-constant- $N_e$  chromosphere
  - mimics Avrett's (1965) [isothermal constant- \$\epsilon\$  two-level-atom scattering atmosphere](#)

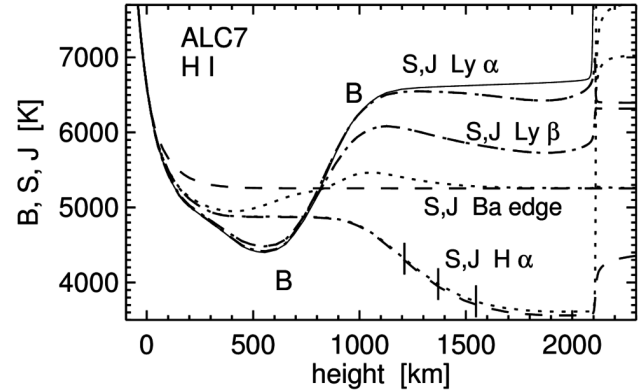
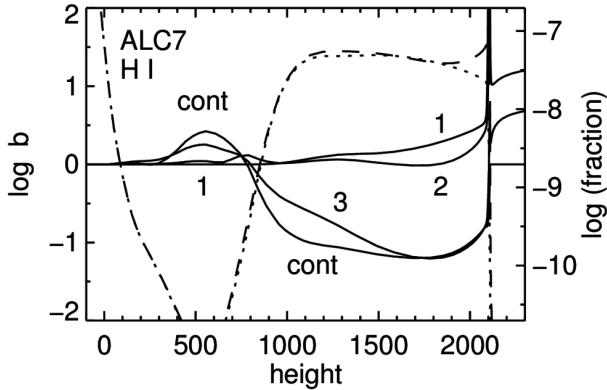
### 3 CRD RESONANT SCATTERING IN AN ISOTHERMAL ATMOSPHERE

RTSA figure 4.12; from Avrett 1965SAOSR.174..101A



- left:  $S/B$  in a plane-parallel isothermal atmosphere with constant  $\epsilon$  for complete redistribution. The curves illustrate the  $\sqrt{\epsilon}$  law and thermalization at  $\Lambda \approx 1/\epsilon$ .
- right: corresponding emergent line profiles and Gaussian extinction profile shape  $\phi$  (only the righthand halves;  $x = \Delta\lambda/\Delta\lambda_D$ )

## 4 HYDROGEN LINES IN THE ALC7 ATMOSPHERE



$H\alpha$ : chromosphere is back-scattering attenuator for radiation from deep photosphere; outward  $S$  decline as in [isothermal constant- \$\epsilon\$  two-level-atom atmosphere](#)

$Ly\alpha$ : tremendous scattering with  $S_{Ly\alpha} \approx J_{Ly\alpha}$  but local thermalization with  $J_{Ly\alpha} \approx B_{Ly\alpha}$  from short photon mean free paths ( $S$  dotted,  $J$  dashed; dot-dashed = identity)

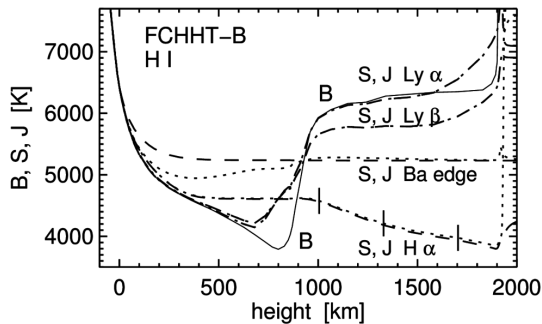
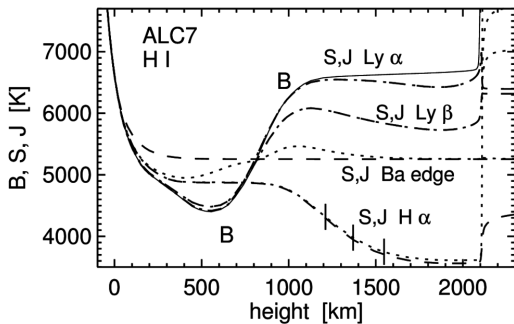
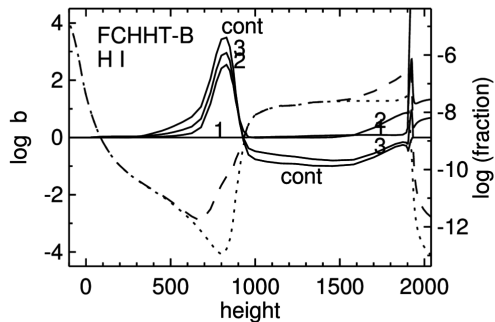
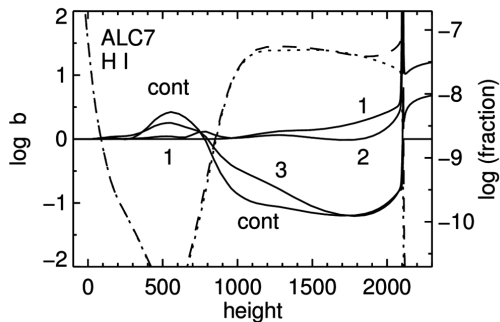
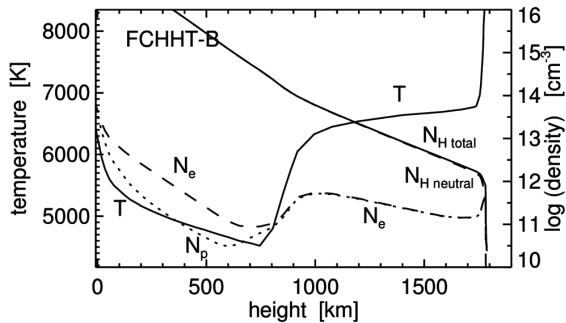
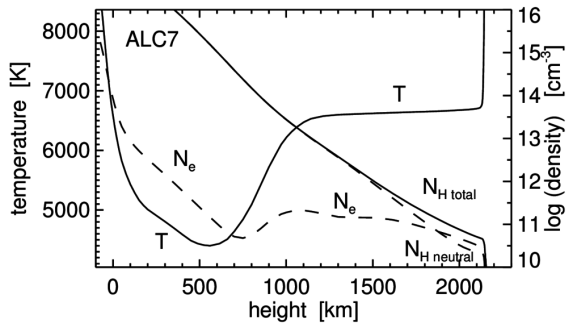
$Ly\beta$ : scattering as  $Ly\alpha$ , shares photon losses in  $H\alpha$  ( $H\alpha$  ticks  $\tau=3, 1, 0.3$ )  
(same  $S/B \approx b_3/b_l$  since  $b_2 \approx b_1$  but offsets differ in temperature representation)

$n=1$ : Saha-Boltzmann  $b_1 \approx 1$  population because hydrogen is neutral  
(except in transition region at right)

$n=2$ : Saha-Boltzmann  $b_2 \approx 1$  population from  $Ly\alpha$  thermalization  
(dotted fraction curve =  $n_2^{LTE}/N_{Htot} \approx$  dashed curve = actual  $n_2/N_{Htot}$ )

ionization:  $b_{cont}/b_2$  defined by SE balancing of  $B(T_{rad}^{Bacont})/B(T_e)$  ionization driving and recombination driving by Balmer and higher line photon losses. The H I top ( $n \geq 2$ ) represents a 3.4 eV alkali atom with ground-state population set by  $Ly\alpha$ .

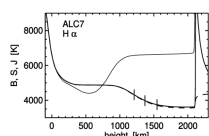
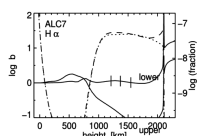
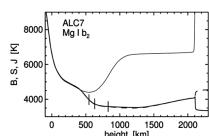
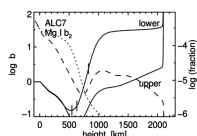
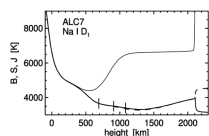
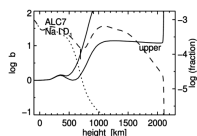
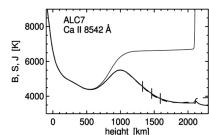
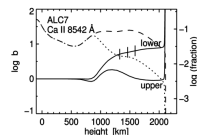
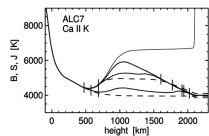
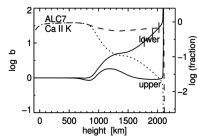
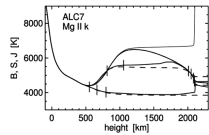
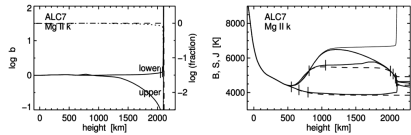
# 5 EXPLAIN EVERYTHING – INCLUDING SIMILARITIES AND DIFFERENCES



## 6 STRONG LINES IN ALC7

Avrett & Loeser 2008ApJS..175..229A

Rutten 2016A&A...590A.124R



### • Mg II k

- extinction LTE, source function 2-level scattering
- high peaks, low PRD dips, low wings

### • Ca II K

- lower abundance and ionization, underionization
- small peaks and PRD dips

### • Ca II 8542

- as Ca II K with Boltzmann lowering and sensitivity
- similar source function sampling as H $\alpha$

### • Na I D<sub>1</sub>

- photospheric scattering, suction and underionization
- no sensitivity to temperature rise

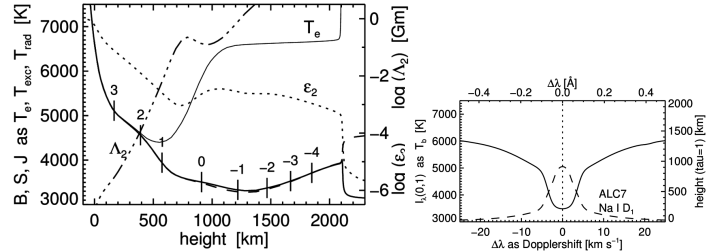
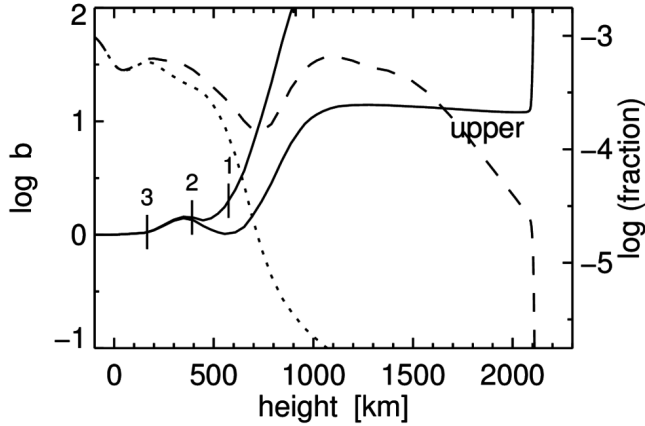
### • Mg I b<sub>2</sub>

- as Na I D<sub>1</sub> but photospheric overionization
- no sensitivity to temperature rise

### • H $\alpha$

- chromospheric scattering of photospheric photons
- chromospheric extinction LTE from Ly $\alpha$  box-up

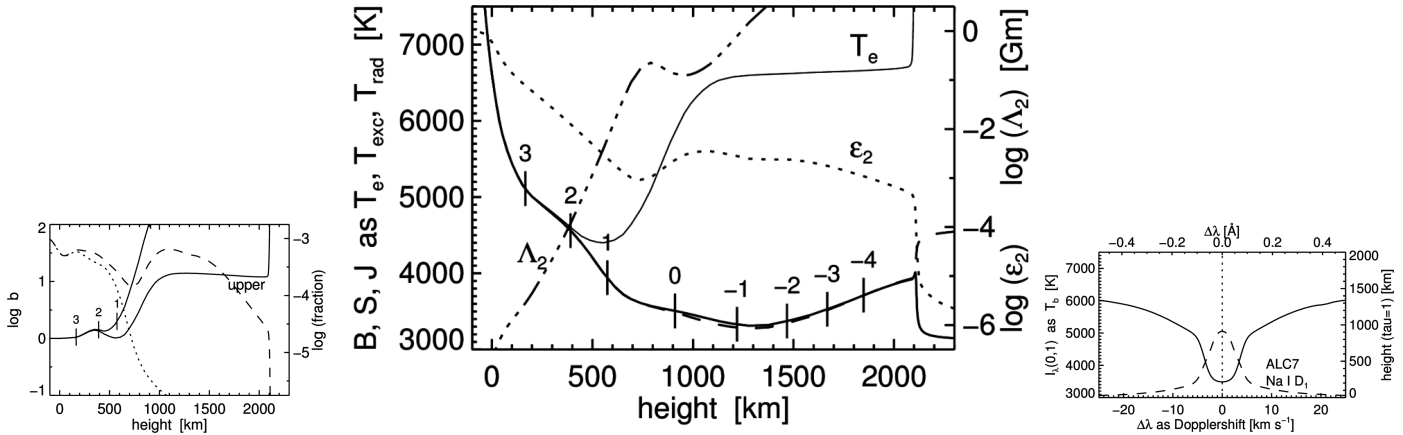
## 7 LINES FROM THE ALC7 ATMOSPHERE: POPULATIONS PLOT (Na I D<sub>1</sub>)



$$b_i \equiv n_i/n_i^{\text{LTE}} \quad \alpha^l \approx b_l \alpha^{\text{LTE}} \quad S^l \approx (b_u/b_l)B \quad \epsilon \approx \epsilon_2 = \alpha^a/(\alpha^s + \alpha^a) \quad S^l \approx (1 - \epsilon_2)\bar{J} + \epsilon_2 B$$

- **solid**: population departure coefficients for Na I D<sub>1</sub>. Unity in deep photosphere from large collision frequency at high density, with  $\epsilon \approx 1$  (*BSJ* plot). Increasing  $b_u < b_l$  divergence =  $S^l < B$  divergence (*BSJ* plot) from  $\sqrt{\epsilon}$ -law resonance scattering. Small initial hump in upper photosphere from photon suction (replenishment from ion reservoir) by scattering-out Na I D photons. Steep  $b_l$  rise above 700 km from ultraviolet underionization ( $1 - c$  edge at 2412 Å, typical for minority neutrals). The  $\log \tau$  ticks on the  $b_l$  curve are for line center.
- **dotted**: fractional population  $n_l^{\text{LTE}}/N_{\text{elem}}$  per Saha-Boltzmann. Scale at right. Na I is a minority species. Initial decrease from increasing ionization at decreasing  $N_e$ , slight hump from less ionization at lower temperature, steep decline at increasing  $T$  and decreasing  $N_e$  (Saha).
- **dashed**: fractional population  $n_l/N_{\text{elem}}$  in NLTE. Line-center optical depth  $\tau_\lambda = -\int(\alpha^l + \alpha^c) dh$  has  $\alpha^l \gg \alpha^c$  and  $\alpha_\lambda^l \sim n_l = (n_l/N_{\text{elem}})A_{\text{elem}}N_{\text{Htot}}$ . Divergence from LTE curve corresponds to departure of  $b_l$  from unity. The steep  $b_l$  increase compensates the steep  $n_l^{\text{LTE}}$  decrease.

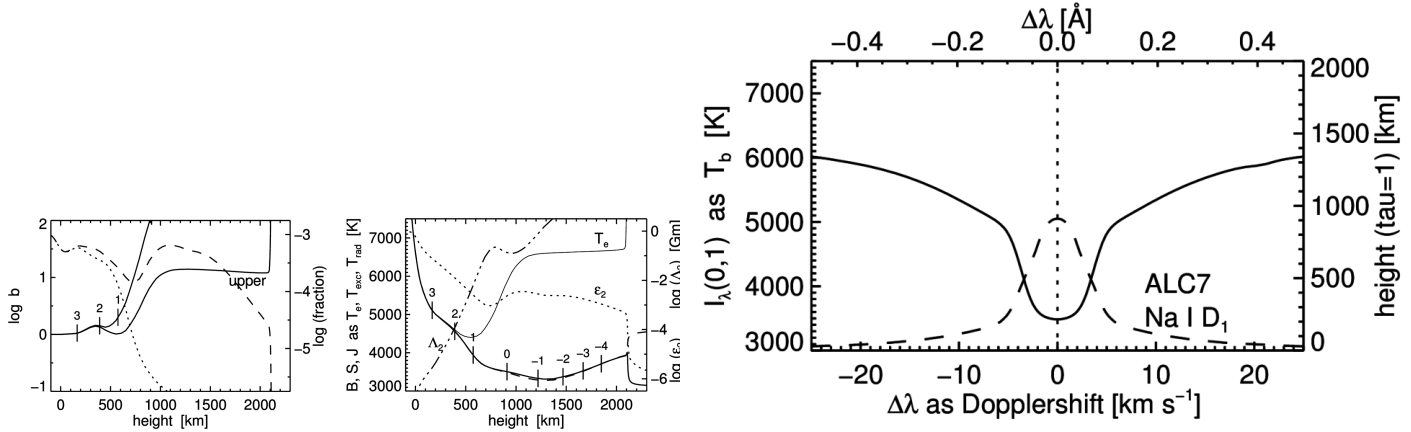
## 8 LINES FROM THE ALC7 ATMOSPHERE: B S J PLOT (Na I D<sub>1</sub>)



$$S^l \approx (b_u/b_l)B \quad \varepsilon \approx \varepsilon_2 \equiv \alpha^a/(\alpha^s + \alpha^a) \quad S^l \approx (1-\varepsilon_2)\bar{J} + \varepsilon_2 B \quad \Lambda_2 \approx \sqrt{\pi}/(\alpha_0 \varepsilon_2)$$

- *thin solid*:  $B_{\lambda_0}$  as temperature  $T_e$  to remove Planck function variation with wavelength for comparison with other lines. The ALC7 atmosphere has a near-isothermal chromosphere.
- *thick solid*: Na I D<sub>1</sub>  $S^l$  as formal excitation temperature  $T_{exc}$ . The  $B > S$  divergence corresponds to the  $b_l > b_u$  divergence in the populations plot, but not equally in their plotted logarithms due to the  $B$  and  $S$  conversions to formal temperature. The  $\log \tau$  ticks are for line center. This scattering line does not sense the ALC7 chromosphere in  $S^l$ .
- *dashed*: profile-averaged angle-averaged intensity  $\bar{J}$  as formal radiation temperature  $T_{rad}$ .
- *dotted*: 2-level photon destruction probability  $\varepsilon_2$  for the Doppler core. Scale to the right. Follows  $N_e$ , so fairly constant over 1000–2000 km from increasing hydrogen ionization.
- *dot-dashed*: 2-level thermalization length  $\Lambda_2$  for the Doppler core in gigameter. Scale to the right. Example:  $\Lambda_2 = -6$  implies thermalization of  $S$  to  $B$  at the center of a 2-km thick feature. The curve label is placed near the line-core thermalization height in the mid photosphere.

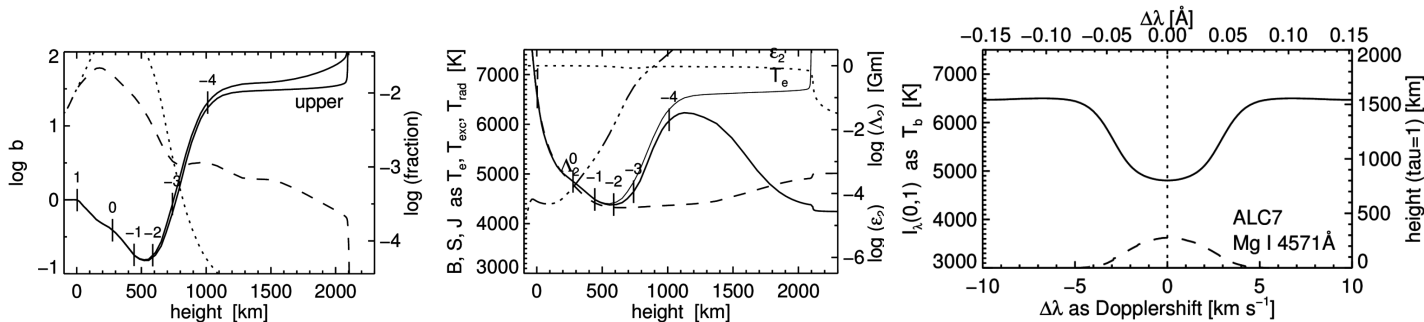
## 9 LINES FROM THE ALC7 ATMOSPHERE: PROFILE PLOT (Na I D<sub>1</sub>)



$$\alpha_\lambda^l \sim b_l (n_l^{\text{LTE}}/N_{\text{elem}}) A_{\text{elem}} N_{\text{Htot}} \varphi(\lambda - \lambda_0) \quad \tau_\lambda = - \int (\alpha^l + \alpha^c) dh \quad I_\lambda(0) \approx S_\lambda^{\text{total}}[\tau_\lambda = 1]$$

- **solid**: emergent intensity in the radial direction, represented as formal brightness temperature for comparison with other lines and the Eddington-Barbier estimate (*BSJ* plot, temperature axes match in the coming line-formation displays). Similarly, the bottom scale for wavelength separation from line center is in km s<sup>-1</sup> for comparison with other lines. Wavelength separations in Å along the top.
- **dashed**:  $\tau_\lambda = 1$  height, scale at right.
- **dotted, vertical**: sampling wavelength(s) for *S* and *J* in the *BSJ* plot. Only one for CRD lines (as Na I D<sub>1</sub>) with frequency-independent line source functions (and  $\bar{J}$  in the *BSJ* plot).
- one might overplot an observed solar disk-center profile, but this is misleading because even a perfect match does not imply that the ALC7 model is correct. ALC7 is an idealized didactic star not like the Sun with an easier-to-understand solar-lookalike spectrum.

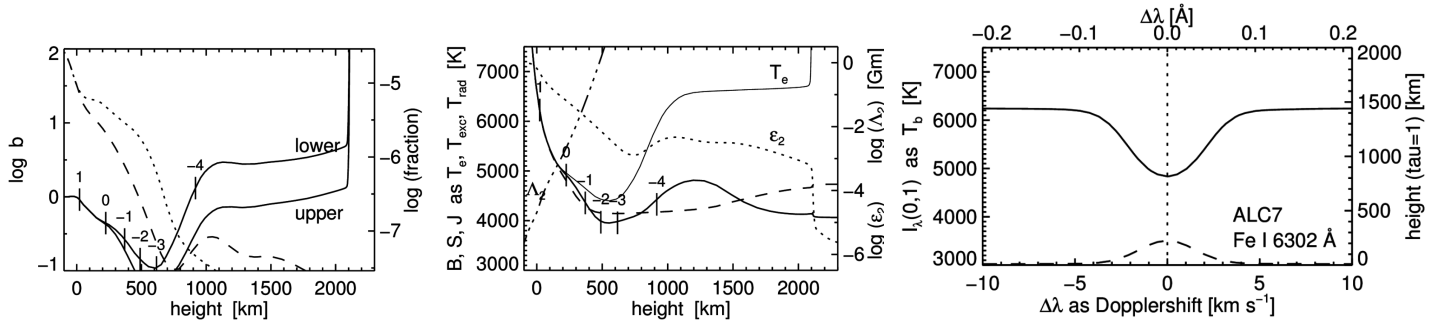
## 10 Mg I 4571 Å FROM THE ALC7 ATMOSPHERE



*unique photospheric line with LTE source function*

- extinction severely out of LTE. Deep  $b_l$  dip across the ALC7 photosphere from overionization by deeply escaping bound-free scattering ultraviolet radiation, including edges of Mg I itself at 2512 and 1621 Å. Corresponding steep  $b_l$  rise above 700 km from ultraviolet underionization where the temperature increases in excess of the ultraviolet radiation temperature.
- this pattern is common to all lines of minority neutrals with ultraviolet ionization wavelengths, including the electron donors (Mg I, Fe I, Si I, Al I).
- source function unusually close to LTE because this is a “forbidden” intersystem line with small  $A_{ul} = 2.7 \cdot 10^2 \text{ s}^{-1}$ , dominated by collisions ( $\epsilon \approx 1$ ) with  $b_u \approx b_l$ ,  $S^l \approx B$  to large heights.
- yet fairly strong because its lower level is the Mg I ground state
- usefulness: photospheric thermometer but requires ultraviolet NLTE for optical depth

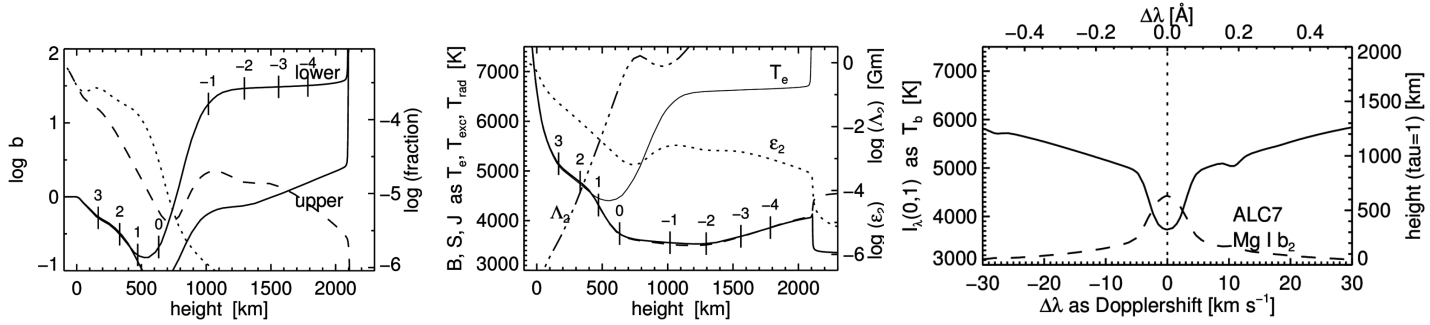
# 11 Fe I 6301.5 Å FROM THE ALC7 ATMOSPHERE



## standard polarimetry line

- severe extinction NLTE across the photosphere due to ultraviolet bound-free scattering overionization and affecting the tau scaling ( $b_l$  curve)
- minor  $S^l$  NLTE from resonance scattering in the upper photosphere ( $S < B$  split)
- “inversion” codes (numerical best-fit iteration) sometimes include  $S^l$  NLTE but usually not extinction NLTE, ignoring that bound-free scattering with  $S^{UV} \approx \bar{J}^{UV}$  depends on 3D temperature gradients in deeper layers and makes  $b_l$  (hence  $n^l$  and  $\alpha^l$ ) non-local both in space and wavelength
- problem: the enormous density of NLTE lines (“haze”) in the ultraviolet affecting  $J^{UV}$
- usefulness: differential line-pair polarimetry with its twin Fe I 6302.5 Å

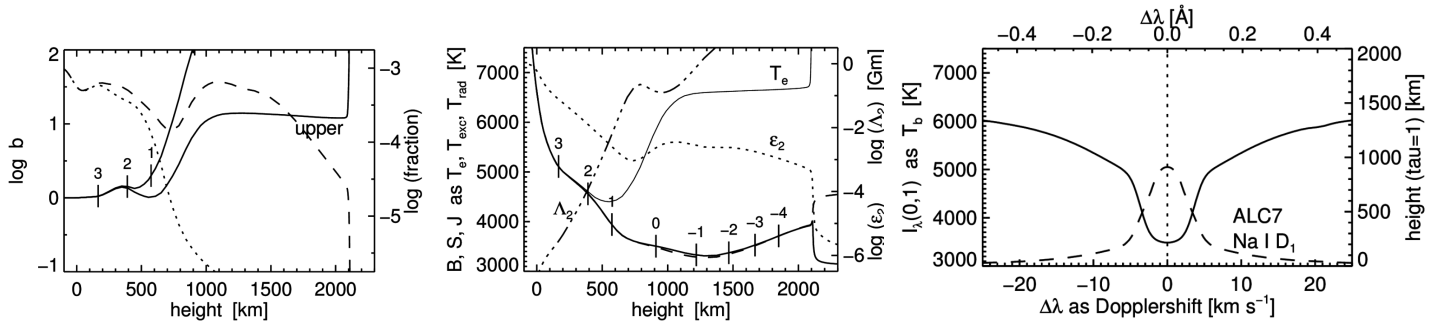
## 12 Mg I b<sub>2</sub> 5173 Å FROM THE ALC7 ATMOSPHERE



*diagnostic of upper photosphere*

- large NLTE  $n_l$  depletion from ultraviolet bound-free scattering across the photosphere
- large NLTE  $b_l$  increase from ultraviolet scattering offsets Saha decline in chromosphere
- CRD scattering source function with  $\epsilon \approx 10^{-3}$
- similar to Na I D<sub>1</sub>
- usefulness: as Na I D<sub>1</sub> but wider core = less asymmetry from reversed granulation

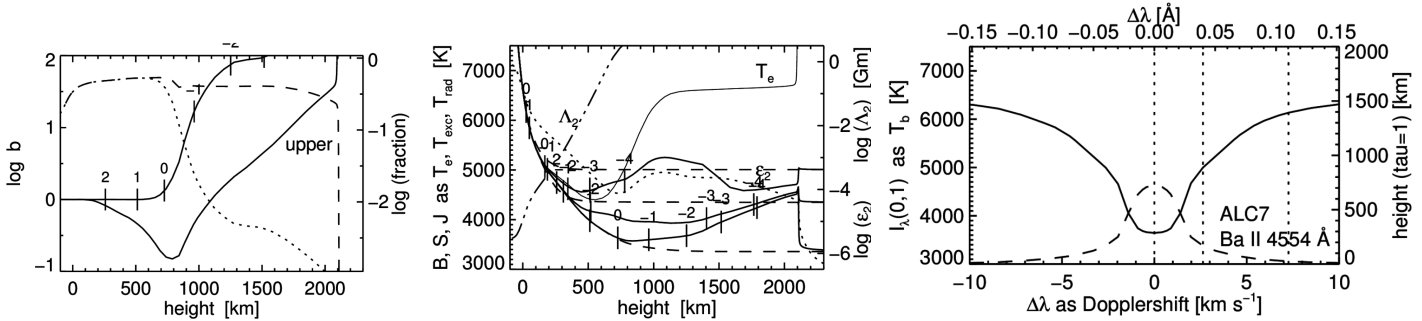
# 13 Na I D<sub>1</sub> 5896 Å FROM THE ALC7 ATMOSPHERE



*Na I D<sub>1</sub>: darkest solar line in optical spectrum = textbook example of two-level scattering*

- photon suction offsets ultraviolet overionization across the photosphere
- ultraviolet underionization offsets Saha depletion above 700 km
- 2-level CRD scattering with  $\varepsilon \approx 10^{-3}$  and  $S \approx \bar{J} \ll B$  in the ALC7 chromosphere
- thermalization in mid photosphere: core intensity does not sense ALC7 chromosphere, observed photons are created near the thermalization depth (height of  $\Lambda_2$  label), observed intensity variation preferentially encodes temperature variation there
- last scattering near  $\tau = 1$ : Doppler and Stokes inner-wing encoding occurs around 500 km
- usefulness: sharp Na I D<sub>1</sub> Dopplergrams indicate deeply-located shocks in fluxtubes

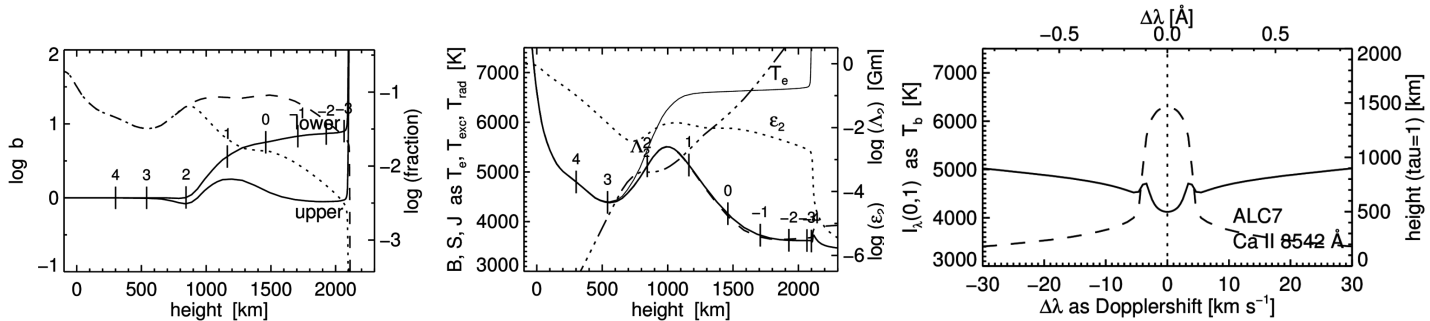
## 14 Ba II 4554 Å FROM THE ALC7 ATMOSPHERE



*weakest PRD line, best velocity diagnostic, good Hanle diagnostic*

- extinction LTE up to line-core formation height thanks to photon losses offsetting overionization (edge at 1240 Å outside Ly $\alpha$ )
- steep  $b_l$  increase above 800 km from underionization offsets Saha depletion
- resonance line with Grotrian diagram similar to Ca II K. PRD scattering source function with  $\epsilon \approx 10^{-4}$ , therefore different monochromatic  $S_\lambda$  and  $J_\lambda$  curves for line center, inner wings, outer wings
- $S^l$  split in upper photosphere produces emission wings at the limb (my 1976 eclipse-expedition PhD thesis)
- usefulness: non-thermal Doppler sensitivity from large mass ( $\sqrt{m_{Ba}/m_H} = 11.7$ ) intricate near-limb Hanle profile from hyperfine structure

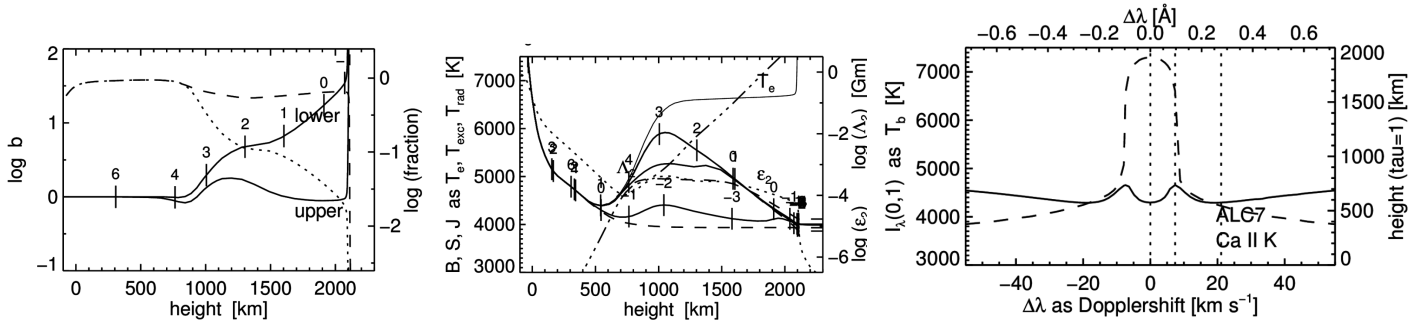
## 15 Ca II 8542 Å FROM THE ALC7 ATMOSPHERE



*cleanest chromospheric diagnostic in the near infrared*

- extinction:  $b_l$  boost from its own photon losses compensates Saha depletion
- CRD scattering source function with  $\epsilon \approx 10^{-2}$
- core formation spans lower ALC7 chromosphere
- best optical line for chromospheric magnetometry
- usefulness: at longer wavelengths more diffraction but less seeing  $\Rightarrow$  prime DKIST line

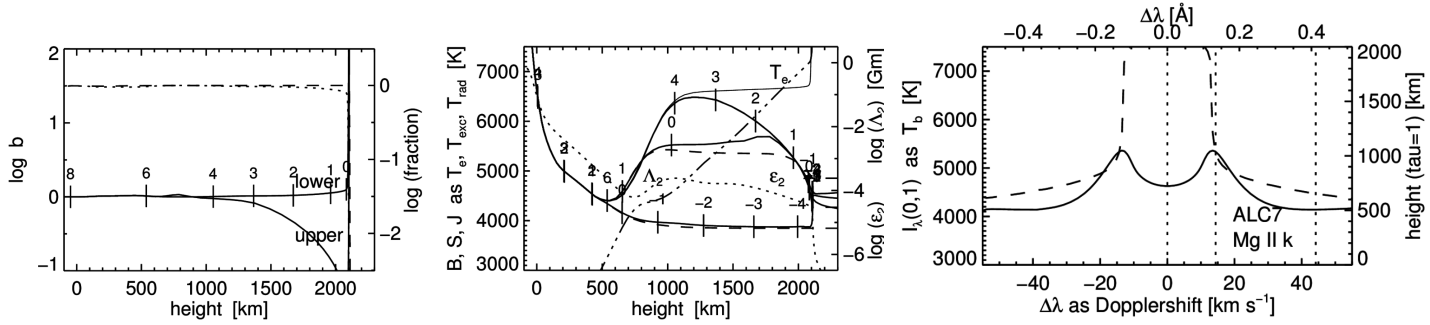
## 16 Ca II K 3934 Å FROM THE ALC7 ATMOSPHERE



*largest extinction in the optical spectrum*

- extinction: successive  $b_l$  boosts from photon losses in infrared triplet and H&K compensates Saha depletion
- PRD scattering source function with  $\epsilon \approx 10^{-4}$  (split between profile center, peaks, dips)
- core formation spans the ALC7 chromosphere
- narrowness of the Doppler core upsets filter imaging so far
- Sunrise-2/SuFi best so far; high hopes for SST/CHROMIS
- usefulness: best optical chromosphere diagnostic but challenging (bandwidth, S/N)

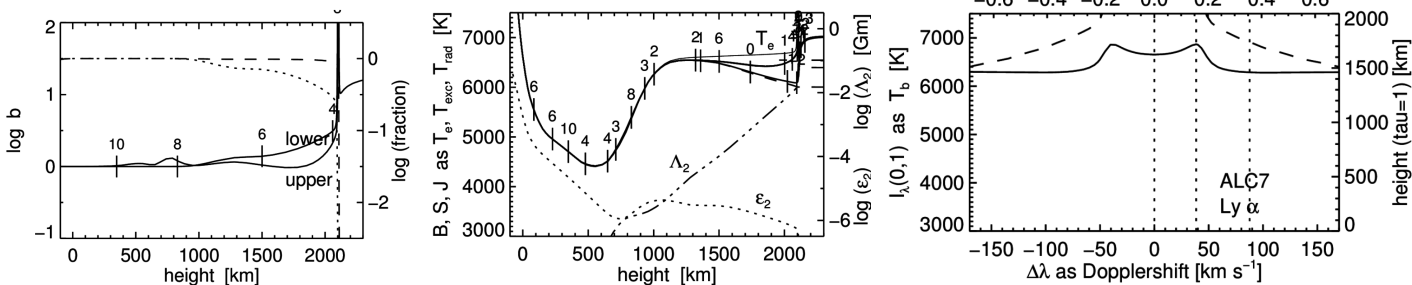
# 17 Mg II k 2796 Å IN THE ALC7 CHROMOSPHERE



*cleanest PRD line and yet larger extinction than Ca II K*

- LTE lower-level population and extinction because all Mg sits in the Mg II ground state
- PRD scattering source function with  $\varepsilon \approx 10^{-4}$  (split between profile center, peaks, dips)
- textbook scattering decline
- similar to Ca II K but with  $18\times$  larger abundance and with much darker wings
- usefulness: key diagnostic but requires space platform slitless imaging spectrometry very difficult combine with “triplet” doublet between h & k (recombination indicator)

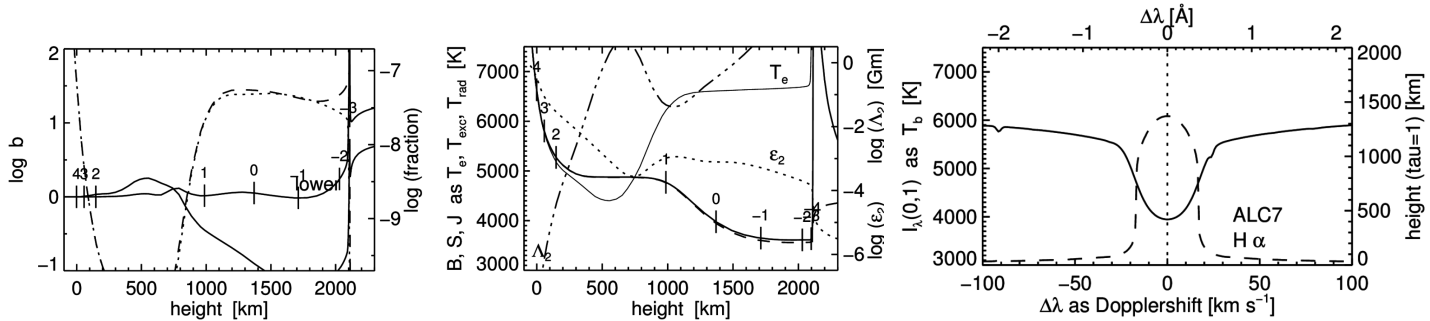
## 18 Ly $\alpha$ 1216 Å IN THE ALC7 CHROMOSPHERE



*champion: largest extinction and most scattering of all lines*

- lower-level population fraction  $\approx 1$ : all hydrogen in ground state
- overpopulation of the ground state towards the transition region from photon losses in wings with slight scattering drops  $S \approx J < B$
- enormous line-center extinction across the ALC7 chromosphere
- PRD scattering source function with  $\epsilon \approx 10^{-6}$  (split between profile center, peaks, dips)
- $\Lambda$  goes from  $\propto 1/\epsilon$  towards  $\propto 1/\epsilon^2$  with density from Stark wing development (not shown)
- radiation lock-in from large extinction produces radiative balance  $n_u(A_{ul} + B_{ul}\bar{J}) = n_l B_{lu}\bar{J}$
- local thermalization from small  $\Lambda$  produces  $\bar{S} \approx B$  throughout ALC7 chromosphere;  $b_u \approx b_l \approx 1$  implies LTE extinction for H $\alpha$  where it escapes
- usefulness: premier diagnostic but needs space; slitless imaging spectrometry difficult

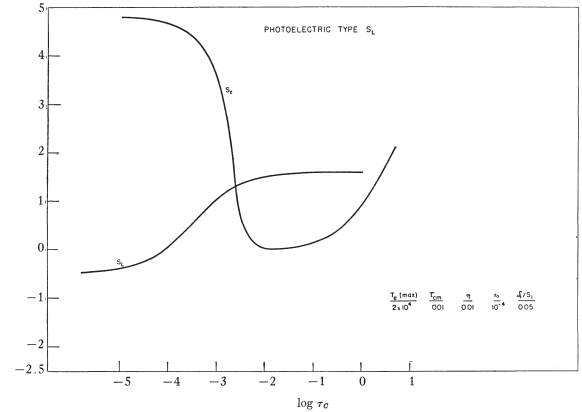
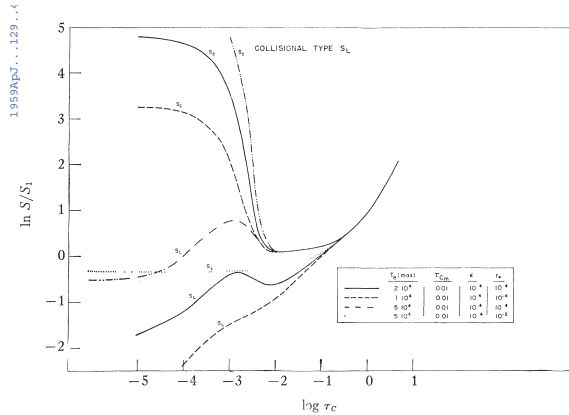
## 19 $H\alpha$ 6563 Å FROM THE ALC7 ATMOSPHERE



$H\alpha$ : extraordinary from high excitation energy, huge element abundance, on top of  $Ly\alpha$

- lower-level fractional population varies  $10^{-10} - 10^{-7}$  due to 10 eV in Boltzmann
- extinction coefficient near-LTE up to 2000 km by  $Ly\alpha$  thermalization
- $S^l \approx$  two-level scattering below transition region (few detours, not “photoelectric”)
- upper photosphere transparent: core shows fibrils, wings show granules
- Eddington-Barbier tau=1 in chromosphere, but photon creation in deep photosphere
- large J across T-min from backscattering: ALC7 chromosphere  $\approx$  scattering attenuator
- wide line core from small atomic mass in Doppler broadening  $\sim \sqrt{2kT/m_H + v_{micro}^2}$
- extended wings from linear Stark effect in deep photosphere (Holtsmark distribution)
- usefulness: prominences, flares, Ellermans, dynamic fibrils, spicules-II, ... = non-E

## 20 CANONICAL CHROMOSPHERIC LINE FORMATION



- CRD line source function including detour paths:

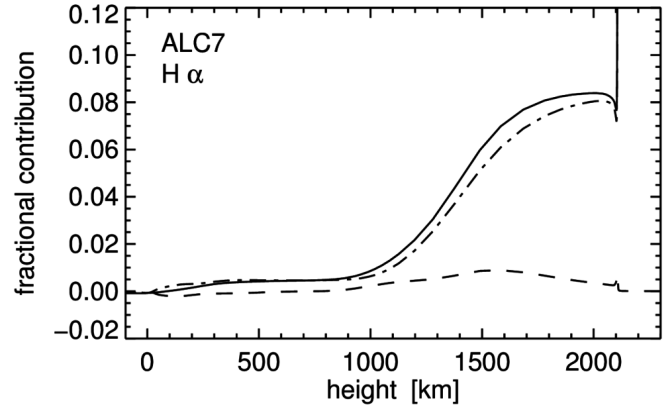
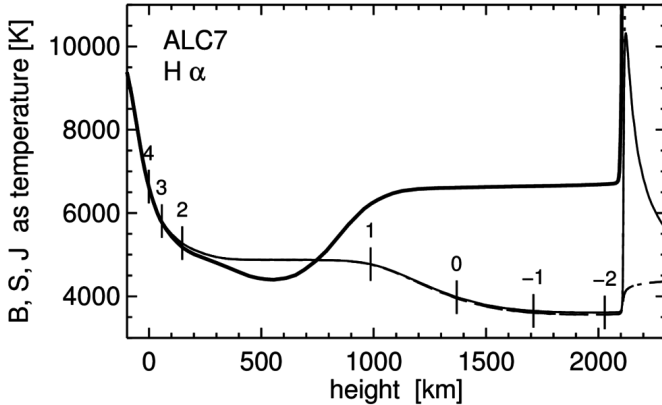
$$S_{\nu_0}^l = \frac{\bar{J}_{\nu_0} + \epsilon'_{\nu_0} B_{\nu_0}(T) + \eta'_{\nu_0} B_{\nu_0}(T_d)}{1 + \epsilon'_{\nu_0} + \eta'_{\nu_0}}$$

$$= (1 - \epsilon_{\nu_0} - \eta_{\nu_0}) \bar{J}_{\nu_0} + \epsilon_{\nu_0} B_{\nu_0}(T) + \eta_{\nu_0} B_{\nu_0}(T_d)$$

- $\epsilon$  = upper-lower collisional destruction fraction of total extinction
- $\eta$  = detour-path extinction fraction of total extinction
- $\epsilon', \eta'$  = idem as ratio to scattering extinction
- $\bar{J}$  = profile-averaged angle-averaged intensity
- $T_d$  = formal detour excitation temperature:  $(g_u D_{ul}) / (g_l D_{lu}) \equiv \exp(h\nu_0 / kT_d)$
- line source function split (Thomas 1957ApJ...125..260T):  
 “collision type” (H & K) or “photoelectric type” ( $H\alpha$ , Balmer continuum feeding)

## 21 H $\alpha$ SOURCE FUNCTION IN THE ALC7 CHROMOSPHERE

after Rutten & Uitenbroek 2012A&A...540A..86R

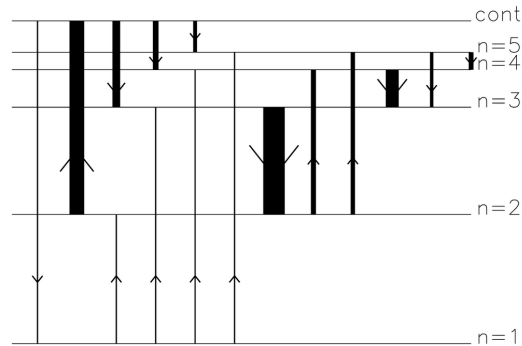


- $$S_{\nu_0}^l = (1 - \varepsilon_{\nu_0} - \eta_{\nu_0}) \bar{J}_{\nu_0} + \varepsilon_{\nu_0} B_{\nu_0}(T) + \eta_{\nu_0} B_{\nu_0}(T_d) = \bar{J}_{\nu_0} + \varepsilon_{\nu_0} [B_{\nu_0}(T) - \bar{J}_{\nu_0}] + \eta_{\nu_0} [B_{\nu_0}(T_d) - \bar{J}_{\nu_0}]$$

The detour part  $\eta_{\nu_0} [B_{\nu_0}(T_d) - \bar{J}_{\nu_0}] / S_{\nu_0}^l$  exceeds the collision part  $\varepsilon_{\nu_0} [B_{\nu_0}(T) - \bar{J}_{\nu_0}] / S_{\nu_0}^l$ . However, their sum  $[S_{\nu_0}^l - \bar{J}_{\nu_0}] / S_{\nu_0}^l$  (solid) reaches only a few percent so  $S_{\nu_0}^l \approx \bar{J}_{\nu_0}$ . Across the ALC7 chromosphere H $\alpha$  is a scattering line, not “photoelectrically controlled”.
- The H $\alpha$  core is dominated by resonance scattering with a formation gap below the chromosphere filled by backscattered radiation. The ALC7 chromosphere acts as scattering attenuator building up its own irradiation. Most emerging photons are created in the deep photosphere where  $\varepsilon_{\nu_0} \approx 1$  and  $\bar{J}_{\nu_0} \approx B_{\nu_0}(T)$ . The granulation pattern has larger contrast than the fibril pattern but is washed out in the scattering across the gap.
- The ALC7 H $\alpha$  core formation is well described by the Eddington-Barbier approximation for an irradiated finite isothermal scattering atmosphere.

## 22 NON-EQUILIBRIUM HYDROGEN IONIZATION IN 1D SHOCKS

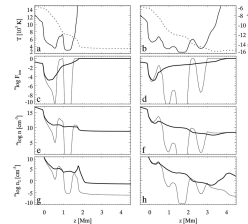
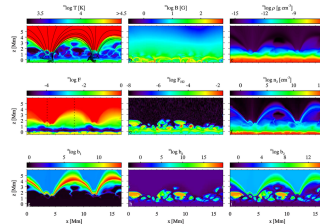
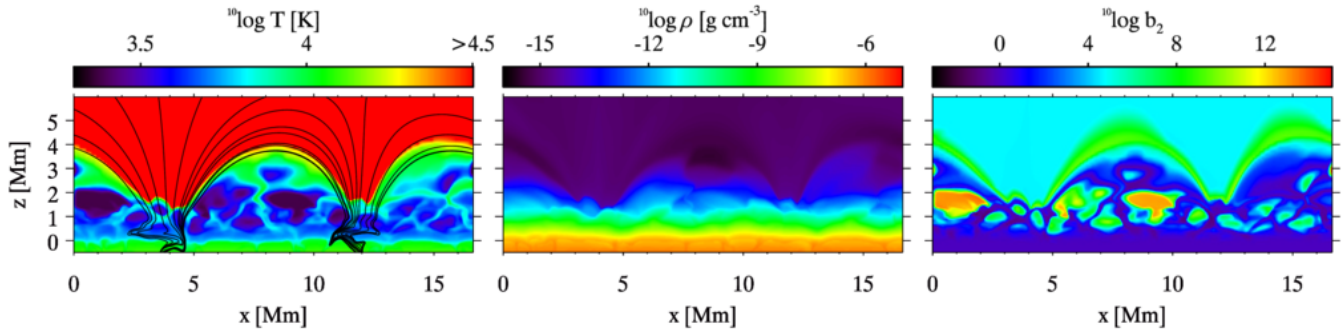
Carlsson & Stein 2002ApJ...572..626C



- RADYN code: 1D(t) hydrodynamics, time-dependent, NLTE radiation, simple PRD
- observed subphotosphere piston drives acoustic waves up that shock near  $h=1000$  km
- Ly $\alpha$  scatters in radiative balance and controls  $n=2$ . Within shocks  $S \approx J$  saturates to  $B$  from radiation lock-in (increased  $\varepsilon$  from partial hydrogen ionization) so that  $b_2 \approx 1$
- collisional Ly $\alpha$  balancing has Boltzmann temperature sensitivity: fast (seconds) in hot gas, slow (minutes) in cool gas, resulting in retardation: post-shock cooling gas maintains the high  $n_2$  shock value at increasing  $b_2$  during minutes, up to huge overpopulation ( $b_2 \approx 10^{10}$ )
- ionization from  $n=2$  in the 3.4 eV alkali-like hydrogen top is an instantaneous statistical-equilibrium balance driven by Balmer continuum  $J \neq B$  and photon losses in Balmer and higher lines, with  $b_{\text{cont}}/b_2 \approx 10^{-1}$  in hot and  $\approx 10^{+3}$  in cool gas, adding to the retarded  $b_2$
- between shocks hydrogen remains hugely overionized versus SE and LTE predictions

## 23 NON-E HYDROGEN IONIZATION IN 2D MHD SHOCKS

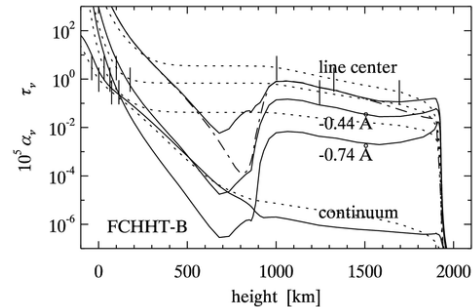
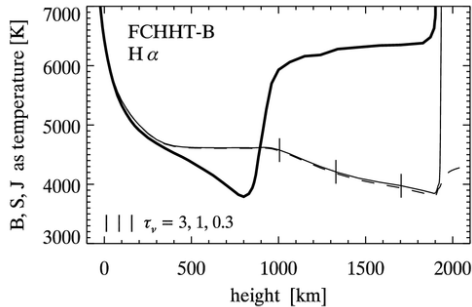
Leenaarts et al. 2007A&A...473..625L



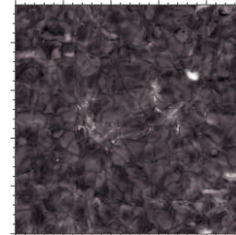
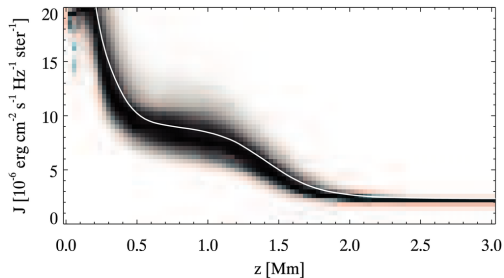
- in shocks  $\text{Ly}\alpha$  has  $S \approx B$  from high  $T$  (fast balancing) and  $N_e$  (10% H ionization)
- retarded collisional balancing in  $\text{Ly}\alpha$ :  $n_2$  hangs near high shock value  $n_2 \approx n_2^{\text{LTE}}$
- gigantic post-shock  $n=2$  overpopulations versus LTE (“S-B underestimates”)
- yet larger post-shock overionization from hydrogen-top Balmer balancing
- approximation: Lyman lines in radiative balance  $\Rightarrow$  green arches artifacts

## 24 H $\alpha$ IN BIFROST

1D plane-parallel: Rutten & Uitenbroek 2012A&A...540A..86R

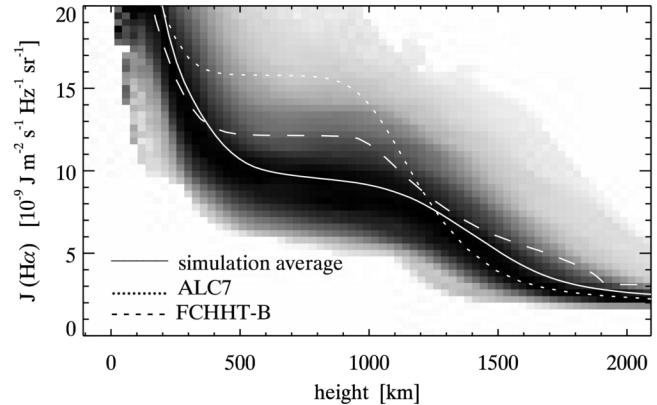
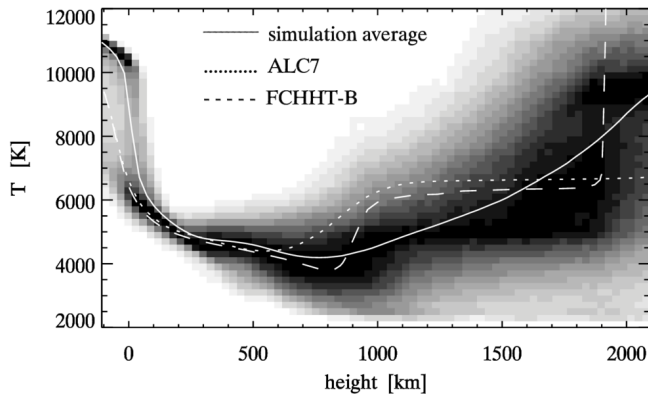


3D non-E MHD snapshot: Leenaarts et al. 2012ApJ...749..136L



- H $\alpha$  is a pure scattering line with  $S \approx J$  and a deep opacity dip in the upper photosphere
- 3D scattering across the opacity gap enhances fibril visibility
- core darkness measures density, core width measures temperature
- caveat: BIFROST snapshot, no non-E line synthesis  $\Rightarrow$  H $\alpha$  severely underestimated

## 25 OSLO SIMULATION VERSUS 1D STANDARD MODELS

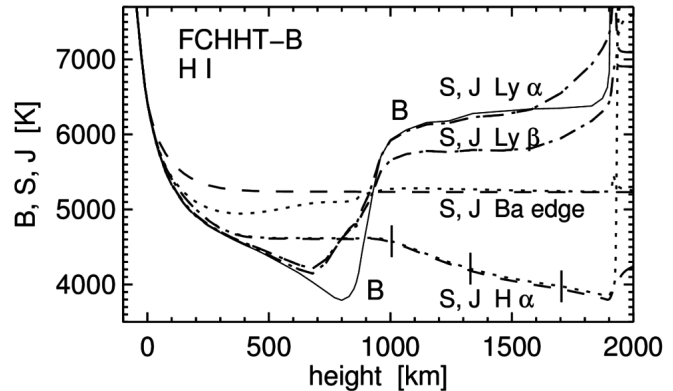
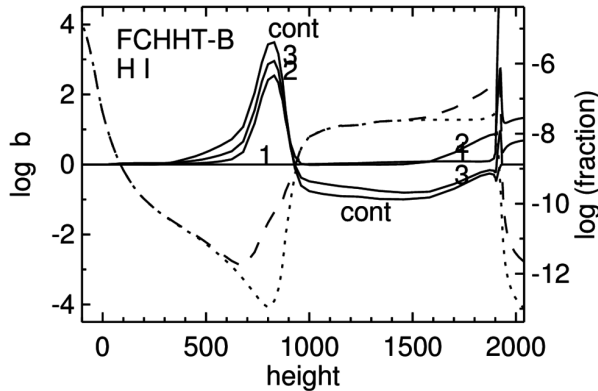


- simulation = state-of-the-art: 3D(t),  $\vec{B}$ , non-HE, SE populations but NE for H  
*Leenaarts, Carlsson & Rouppe van der Voort* [2012ApJ...749..136L](#)
- ALC7 = UV fit: 1D static, no  $\vec{B}$ , HE + microturbulence, SE populations  
*Avrett & Loeser* [2008ApJS..175..229A](#)
- FCHHT-B = UV fit: 1D static, no  $\vec{B}$ , HE + imposed acceleration, SE populations  
*Fontenla, Curdt, Haberreiter, Harder & Tian* [2009ApJ...707..482F](#)

The  $T$  and  $J_\nu(\text{H}\alpha)$  behavior seems arguably similar. However, the conceptual differences between plane-parallel static hydrostatic-equilibrium modeling and the 3D(t) MHD simulation are enormous (cf. Newtonian gravitation versus general relativity). The  $T(h)$  stratifications in the simulation vary tremendously, with shocks propagating upwards and sideways and the increase to coronal temperature dancing up and down over a large height range.

## 26 HYDROGEN AUREOLE BOOSTING IN COOL GAS BESIDE HOT GAS

Fontenla et al. 2009ApJ...707..482F Rutten & Uitenbroek 2012A&A...540A..86R Rutten 2016A&A...590A.124R



FCHHT-B: steep  $B$  rises to chromosphere and corona emulate adjacent cool and hot features

$\text{Ly}\alpha$ : scattering back-radiation boosts  $S_{\text{Ly}\alpha} \approx J_{\text{Ly}\alpha}$  and  $\text{H}\alpha$  extinction  $\propto b_2 \approx S_{\text{Ly}\alpha}/B_{\text{Ly}\alpha}$  towards hot-feature value (left: dotted  $n_2^{\text{LTE}}/N_{\text{Htot}}$ , dashed actual  $n_2/N_{\text{Htot}}$ )

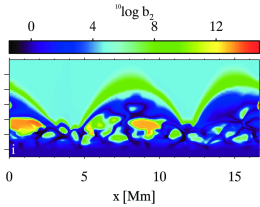
ionization: also  $b_2$ -boosted, with additional  $b_{\text{cont}}/b_2$  offset defined by the Balmer continuum

$\text{Ly}\beta$ :  $b_3$  between  $b_2$  and  $b_{\text{cont}}$  and sharing  $\text{H}\alpha$  photon losses

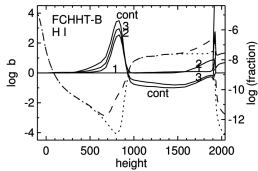
$\text{H}\alpha$ : the FCHHT-B chromosphere is a back-scattering attenuator just as in the [ALC7 atmosphere](#). The  $b_2$  peak from  $\text{Ly}\alpha$  irradiation does not affect  $\text{H}\alpha$  because even with this boost the  $\text{H}\alpha$  extinction in the temperature minimum remains negligible.

A hot feature embedded in cooler gas has a similar  $\text{Ly}\alpha$  scattering aureole enhancing H ionization and  $\text{H}\alpha$  extinction around it. A temporary hot disturbance leaves such spread-out boost behind (a wake when moving).

## 27 H $\alpha$ EXTINCTION RECIPE



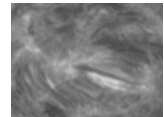
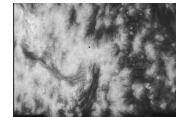
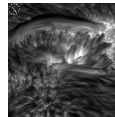
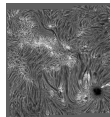
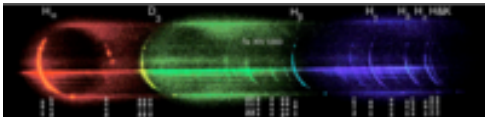
- *retarded Ly $\alpha$  balancing: extinction memory of hot moments*
  - $n_2 \approx n_2^{\text{LTE}}$  in hot shocks from fast Ly $\alpha$  balancing and increased  $\varepsilon$
  - $n_2$  decays slowly, tracking high shock values
  - gigantic  $b_2$  in post-shock cooling clouds until next shock



- *Ly $\alpha$  scattering: aureole boosting*
  - Ly $\alpha$  scattering defines  $S_{\text{Ly}} \approx J_{\text{Ly}}$  with radiative balance
  - hot features in cool gas have Ly $\alpha$  scattering aureoles
  - H I top ( $n \geq 2$  including  $n_{\text{ion}}$ ) boosted in aureoles

- *H $\alpha$  extinction recipe*

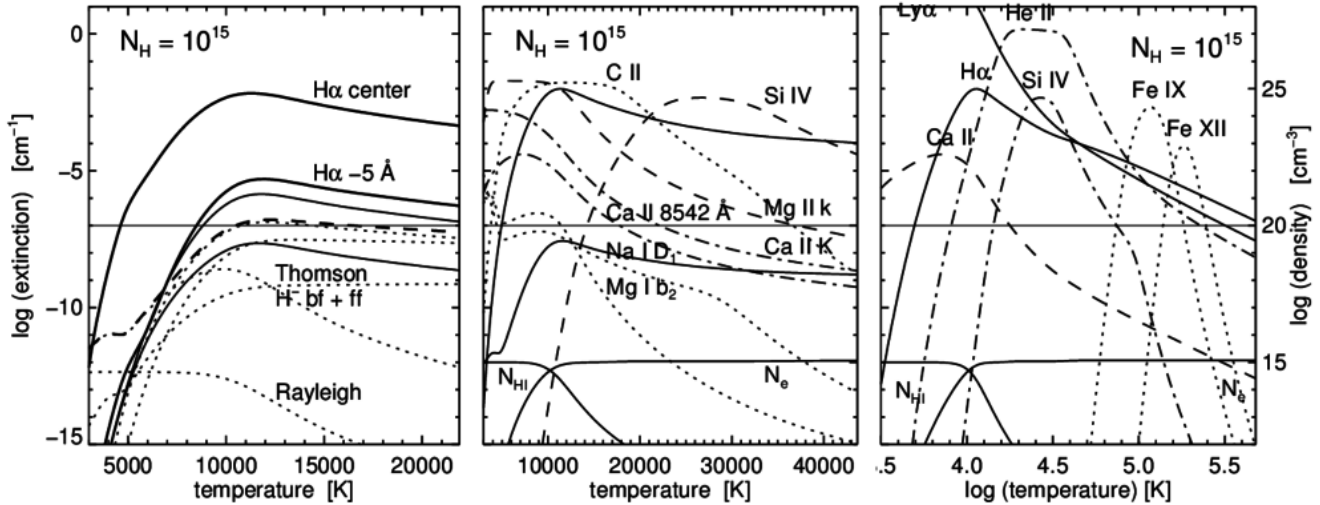
- find hottest instance nearby ( $\sim 300$  km) and in recent past ( $\sim$ minutes)
- compute Saha-Boltzmann fractional  $n = 2$  population then and there
- use this extinction value in cooler gas around it and afterwards
- small hot features leave wider H $\alpha$  marks (as the grin of the Cheshire cat)
- fast small hot features leave wider H $\alpha$  trails (as contrails from jet engines)



## 28 $H\alpha$ HOT-ONSET EXTINCTION

Rutten 2016A&A...590A.124R

Rutten & Rouppe van der Voort 2016arXiv160907616R



- bachelor exercise “Cecilia Payne”: compute  $N_e$  and  $\kappa^{LTE}$  for given  $N_H$
- line at  $\log(\alpha) = -7$ :  $\tau = 1$  across 100-km slab
- $H\alpha$  moustaches beyond continuum: linear Stark + Holtmark
- photosphere:  $H\alpha$  much weaker than Ca II H & K, Ly $\alpha$  mm photon paths
- chromosphere:  $H\alpha \approx$  Ca II 8542 Å, opposite temperature sensitivity
- clapotispheric shocks:  $H\alpha \sim$  Ca II H & K and remembers that
- outburst onsets: extraordinary  $H\alpha$  opacity, competing with (E)UV lines